

Orbital Estimation of Binary Stars

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1 cent piece of advice

Become an expert in something nobody is willing to learn but a lot of people need (e.g. data processing and handling, especially of binaries and other strange objects, for a space astrometry mission).

Overview

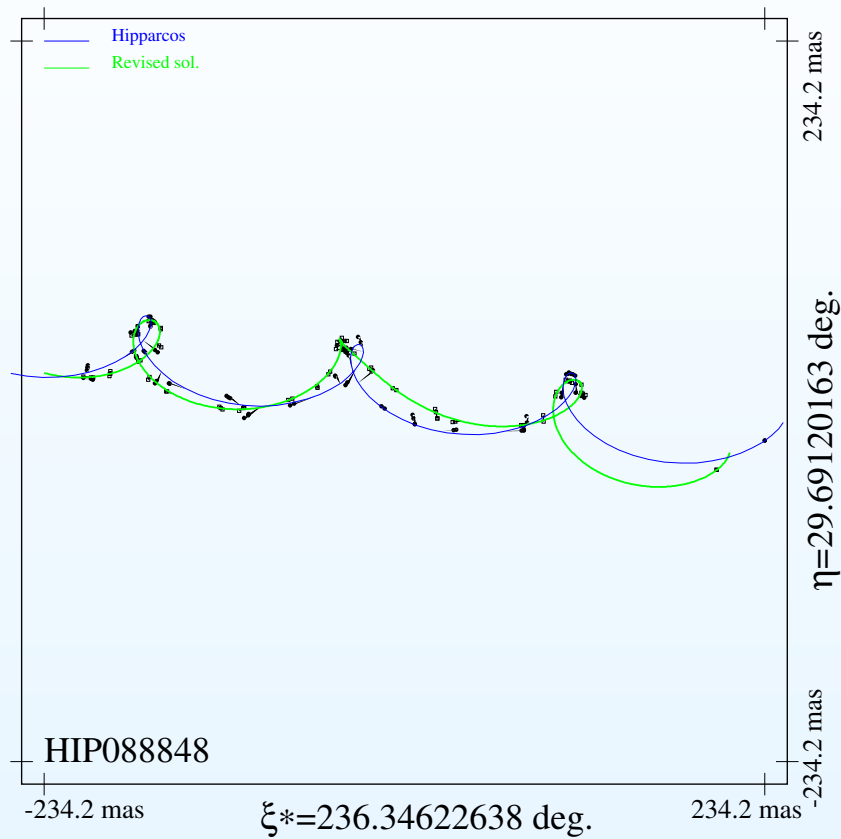
1. Visual, astrometric binaries and extrasolar astrometric orbit fitting

- What does one see and how to model it?
- Two-body problem
- Why does one care? Stellar and planetary systems quite similar
- How to guess the orbit?
- How to efficiently fit it?
- 3+ stellar and planetary systems not so similar after all
- Preliminary Conclusions

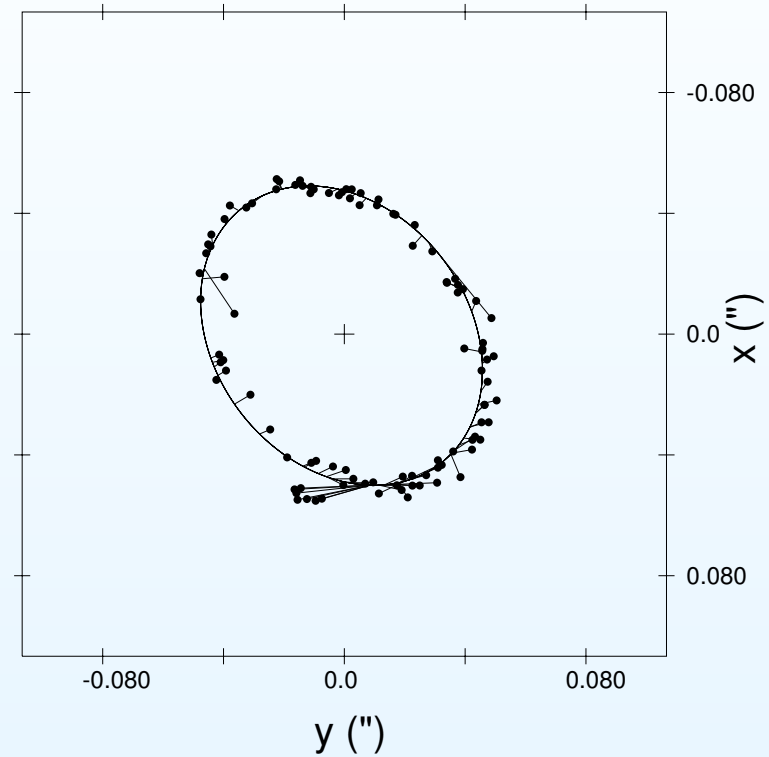
2. Spectroscopic watch dog

3. Blind fit

Absolute path versus relative orbit

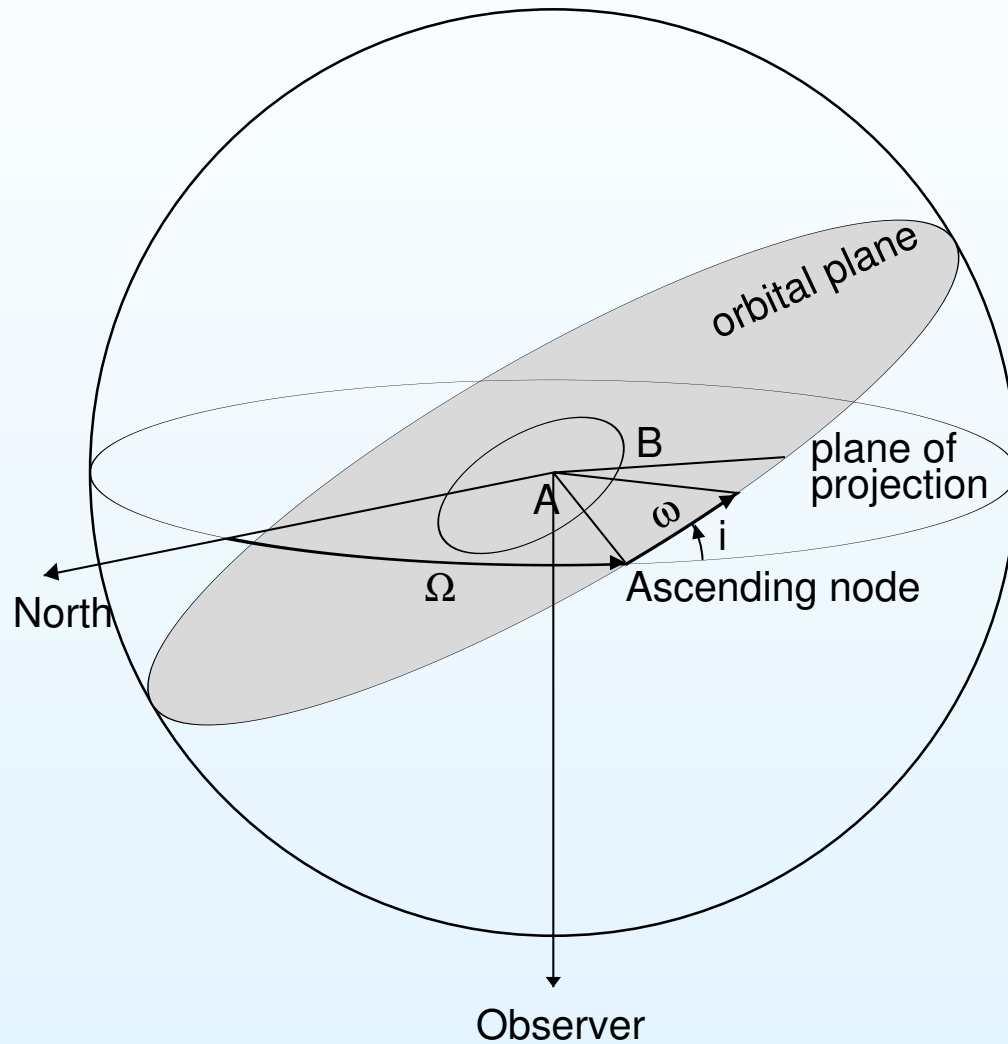


05167+4600Aa



The relative visual orbit requires 7 parameters only.

Astrometric or visual binaries and extrasolar planets



- a : semi-major axis of the orbit
- e : eccentricity
- P : orbital period
- T : epoch of the periastron passage
- i : inclination
- ω : argument of the periastron
- Ω : longitude of the ascending node

Two-body absolute motion

$$\xi = \alpha_0^* + P_\alpha \varpi + (t - t_0) \mu_{\alpha^*} + B(\cos E - e) + G\sqrt{1 - e^2} \sin E,$$

$$\eta = \delta_0 + P_\delta \varpi + (t - t_0) \mu_\delta + A(\cos E - e) + F\sqrt{1 - e^2} \sin E,$$

$$E = \frac{2P}{\pi}(t - T_0) + e \sin E,$$

$$\alpha^* \equiv \alpha \cos \delta,$$

$$A \equiv a(\cos \omega_1 \cos \Omega - \sin \omega_1 \sin \Omega \cos i),$$

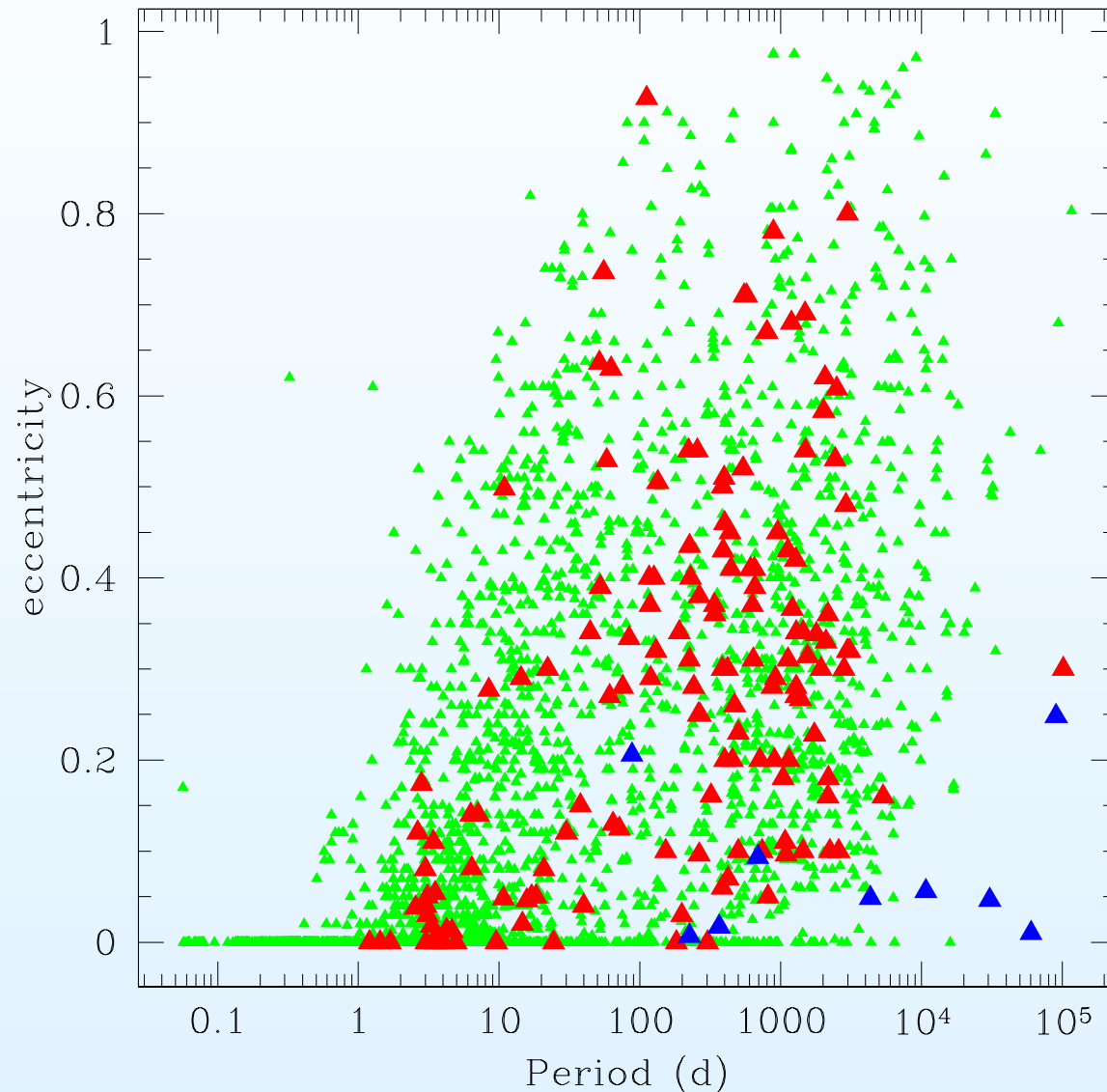
$$B \equiv a(\cos \omega_1 \sin \Omega + \sin \omega_1 \cos \Omega \cos i),$$

$$F \equiv a(-\sin \omega_1 \cos \Omega - \cos \omega_1 \sin \Omega \cos i),$$

$$G \equiv a(-\sin \omega_1 \sin \Omega + \cos \omega_1 \cos \Omega \cos i).$$

12-parameter nonlinear model

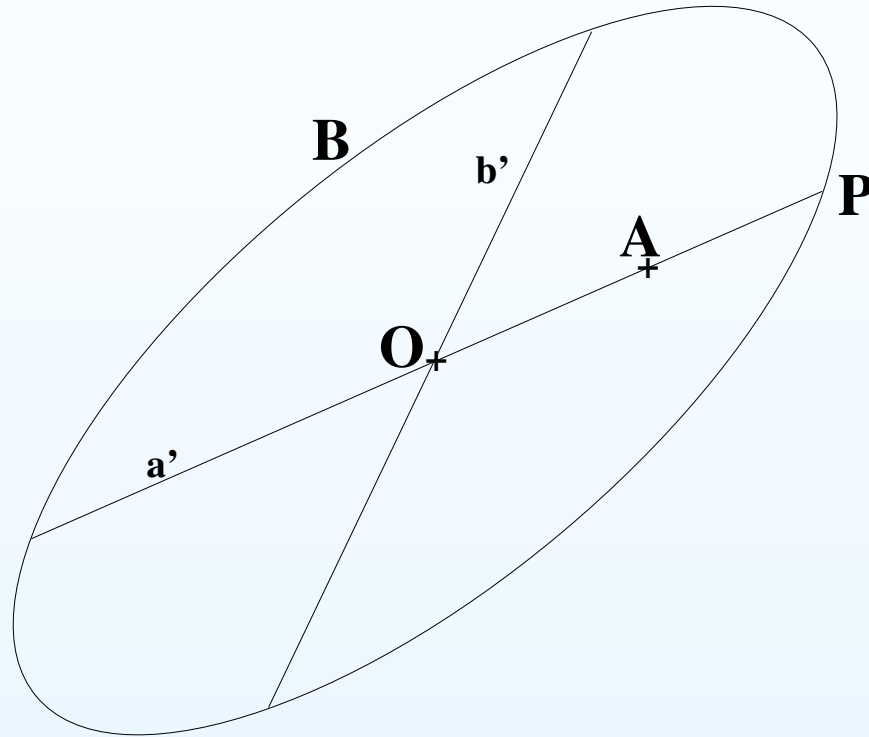
Spectroscopy-based e-log P



No clear difference between stellar and planetary systems.

Codes written for ones can be used for the others.

Old goodies: geometric method – Zwiers' method



Orbital solution derived from the apparent ellipse.

Instead of drawing it, fit

$$Ax^2 + 2Hxy + By^2 + 2Gx + 2Fy + J = 0.$$

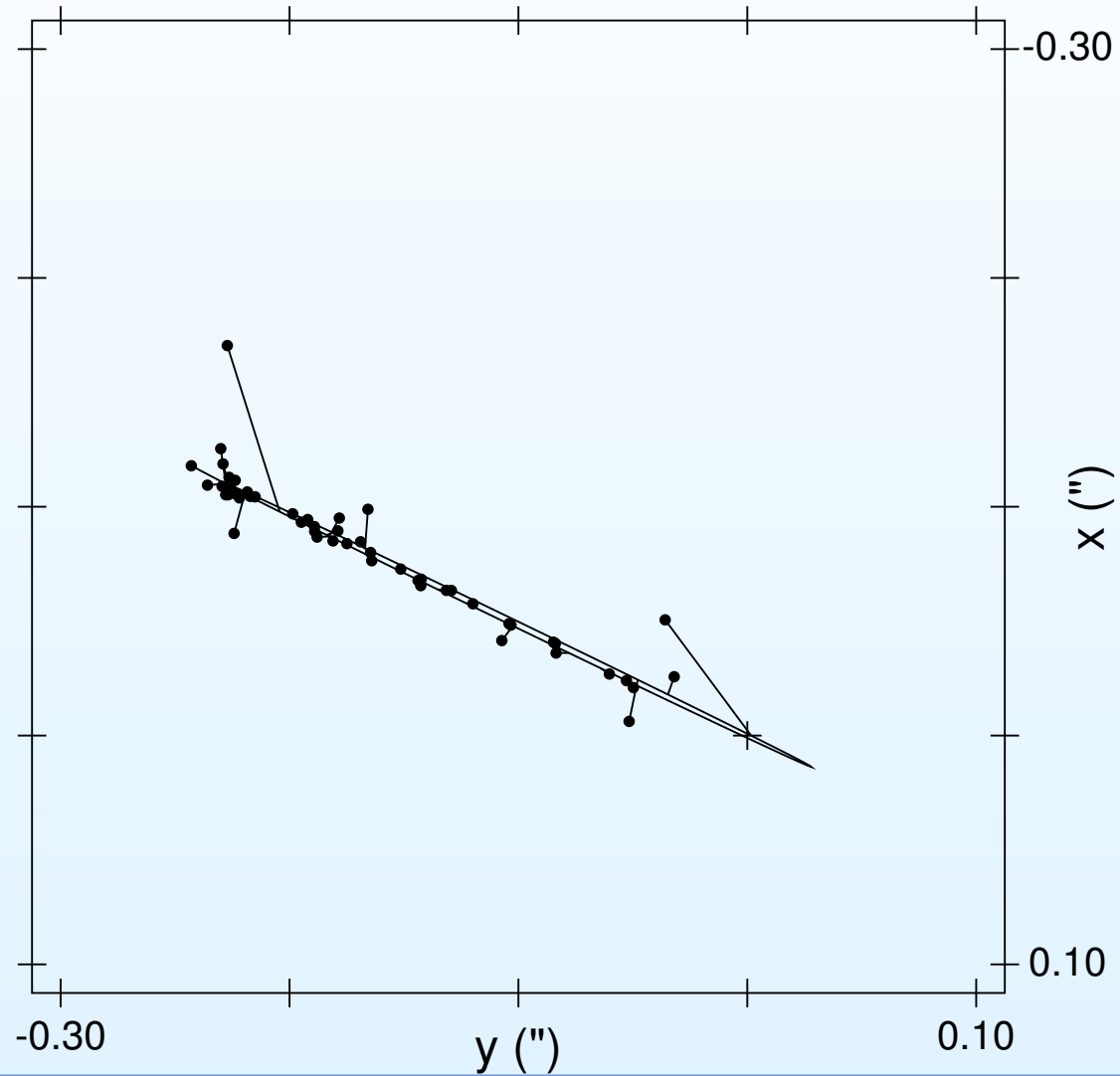
Avoid assuming $J = 1$ (numerical stability). The constants are **not** the Thiele-Innes ones.

Weaknesses:

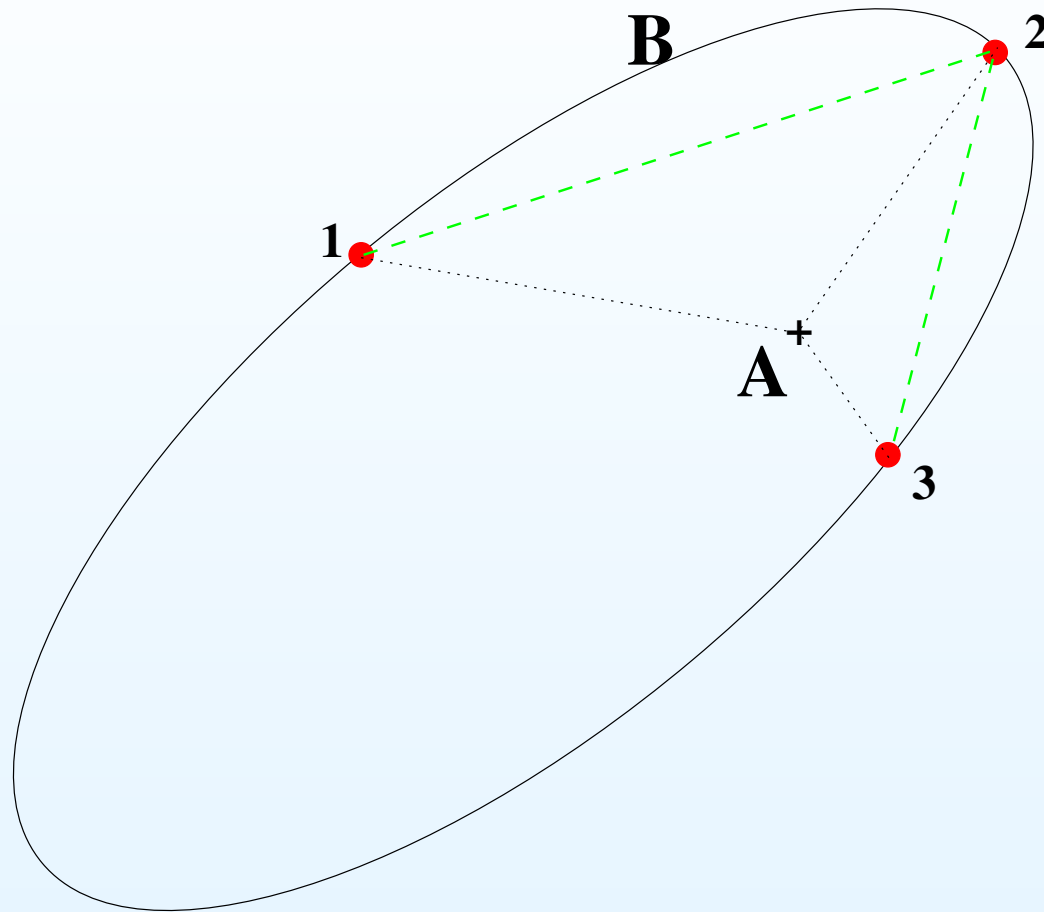
- The apparent orbit needs to be drawn.
- Cannot be applied when $i = 90^\circ$ (degeneracy).

Independent term

HIP 14328



Old goodies: Thiele-Innes-Van Den Bos' method



Exact solution of the equations for 3 *normal* points and the projected constant of areas ($\rho^2 \frac{d\theta}{dt}$). A better be in triangle(1,2,3).

Alternative: replacing $\rho^2 \frac{d\theta}{dt}$ with a 4th point.

High sensitivity of the overall solution on those favored 3 or 4 points.

What if 1D-observations?

Fourier

- Keplerian orbit = periodic phenomenon.
- Fourier expansion of the positions (Monet 1979, ApJ, 234, 275).

$$O(t_i) = \sum_{n=0}^{\infty} a_n \cos(n\mathcal{M}_i) + \sum_{n=1}^{\infty} b_n \sin(n\mathcal{M}_i)$$

$$\mathcal{M}_i = \frac{2\pi}{P}(t_i - t_0)$$

$$O(t_i) = \sum_{n=0}^{\infty} \alpha_n \cos(nM_i) + \sum_{n=1}^{\infty} \beta_n \sin(nM_i)$$

$$\Delta = \frac{2\pi}{P}(T - t_0)$$

$$\begin{aligned} M_i &= \frac{2\pi}{P}(t_i - T) \\ &= \mathcal{M}_i - \Delta \end{aligned}$$

Fourier (cnt.)

- Reverse problem: deriving the orbital parameters from the Fourier coefficients (P assumed).

$$\sum_{n=0}^{\infty} \alpha_n \cos(nM_i) + \sum_{n=1}^{\infty} \beta_n \sin(nM_i) = f(o.e.) \sum_{n=0}^{\infty} F_n(e) \cos(nM_i) + g(o.e.) \sum_{n=1}^{\infty} G_n(e) \sin(nM_i)$$

- Tradeoff between solving the minimum number equations exactly or least squares on many equations. Usually the former.
- Can be applied to incomplete observations (e.g. ρ only).
- The lower the eccentricity the better.

Global minimization

Fitting the nonlinear model directly (rather a linear expansion) using

- Simulated Annealing,
 - Boltzmann probability, temperature.
 - Point generator.
- Genetic Algorithm
 - Encoding, elitism, crossover, mutations
 - Parallel computing
- Tabu search,
 - Diversification, intensification, tabu list, . . .

No guaranteed convergence to the global minimizer within a finite time.

Could be applied to a subset of parameters only.

A two-step approach

In space astrometry, single star solution is always derived anyway (GIS):

$$\Xi^2 = \left(\Delta v - \sum_p \frac{\partial v}{\partial p} \Delta p - \frac{\partial v}{\partial o} o \right)^t \mathbf{V}^{-1} \left(\Delta v - \sum_p \frac{\partial v}{\partial p} \Delta p - \frac{\partial v}{\partial o} o \right)$$

where

$$\frac{\partial v}{\partial o} = \frac{\partial v}{\partial \alpha_0^*} \frac{\partial \xi}{\partial o} + \frac{\partial v}{\partial \delta_0} \frac{\partial \eta}{\partial o}$$

Problem: $O(e^{12})$ local minima. A good initial guess of the solution is required.

$\forall e, P \ \& \ T, \exists!$ minimizer of $\Xi_{12}^2(\mathbf{p}, A, B, F, G, e, P, T)$

Instead $\Xi_3^2(e, P, T)$ min! where

$$\Xi_3^2(e, P, T) = \min \Xi_{12}^2(\mathbf{p}, A, B, F, G, e, P, T)$$

Speed up!

For any matrix symmetric positive definite matrix M , there is an upper triangular matrix R such that $M = R^t R$ (Cholesky decomposition).

So, instead of evaluating Ξ^2 using V^{-1} , find R such that $V^{-1} = R^t R$ and then

$$\begin{aligned}\Xi^2 &= \left(\Delta v - \sum_p \frac{\partial v}{\partial p} \Delta p - \frac{\partial v}{\partial o} o\right)^t V^{-1} \left(\Delta v - \sum_p \frac{\partial v}{\partial p} \Delta p - \frac{\partial v}{\partial o} o\right) \\ &= \left(\Delta v - \sum_p \frac{\partial v}{\partial p} \Delta p - \frac{\partial v}{\partial o} o\right)^t R^t R \left(\Delta v - \sum_p \frac{\partial v}{\partial p} \Delta p - \frac{\partial v}{\partial o} o\right) \\ &= \left(R \Delta v - \sum_p R \frac{\partial v}{\partial p} \Delta p - R \frac{\partial v}{\partial o} o\right)^t \left(R \Delta v - \sum_p R \frac{\partial v}{\partial p} \Delta p - R \frac{\partial v}{\partial o} o\right)\end{aligned}$$

Each iteration becomes $O(N)$ instead of $O(N^2)$

3+ component model

Basic gravitational interaction, no tidal effect, no mass transfer, ...

$$\begin{aligned}\xi &= \xi_S + \sum_k \left(B_k (\cos E_k - e_k) + G_k \sqrt{1 - e_k^2} \sin E_k \right), \\ \eta &= \eta_S + \sum_k \left(A_k (\cos E_k - e_k) + F_k \sqrt{1 - e_k^2} \sin E_k \right), \\ E_k &= \frac{2P_k}{\pi} (t - T_{0,k}) + e_k \sin E_k,\end{aligned}$$

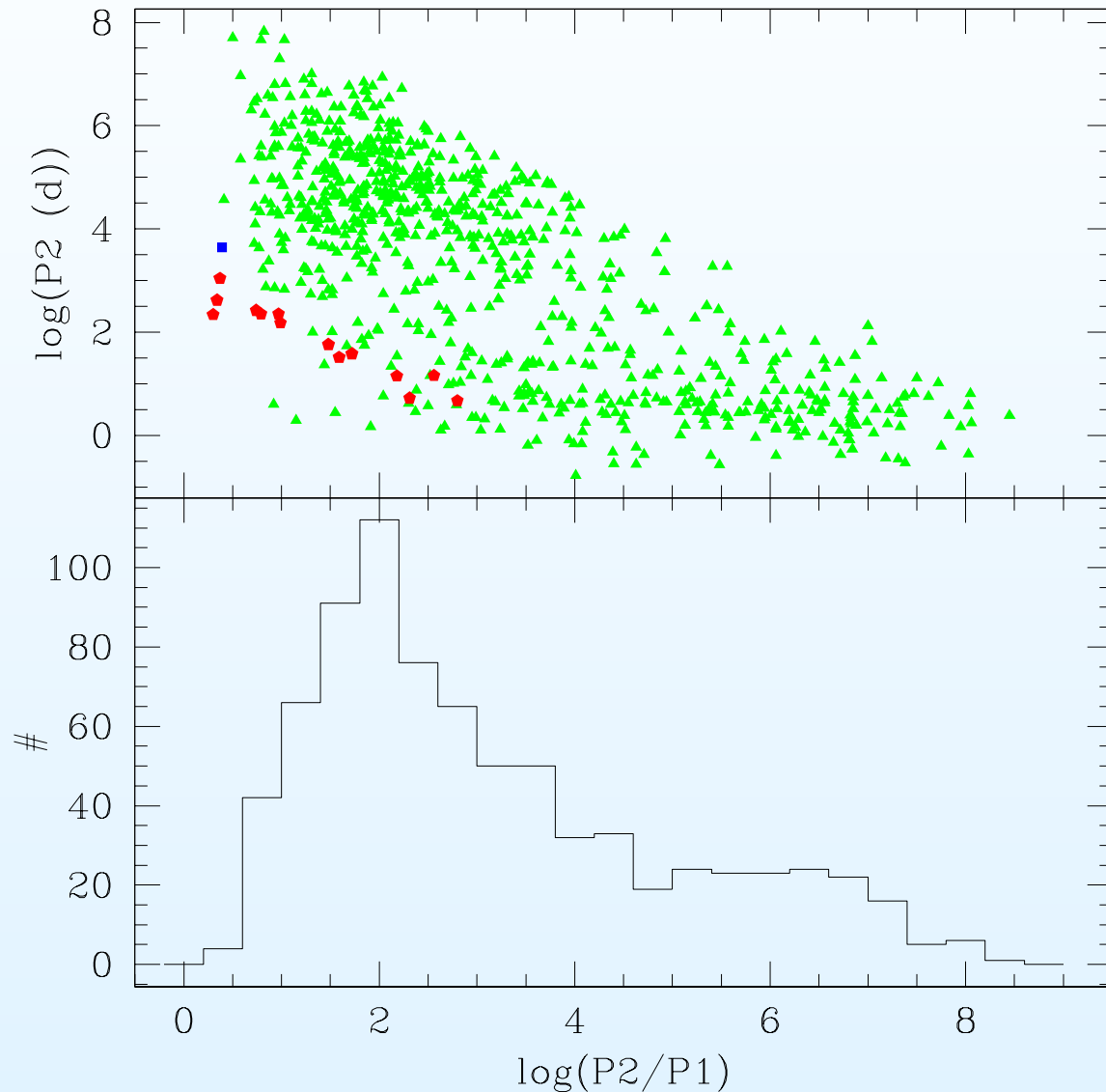
5+7*N-parameter nonlinear model.

Period ratio

Tokovinin (1997)
Scheinerd's enc.

Planetary periods
are more similar
than stellar ones.

No longer hierar-
chical ($3 \neq 2+2$).



Conclusions

- Different interpretations but equations are essentially identical for visual and astrometric binary and extrasolar planetary orbits.
- Unique code for the two families.
- Fitting 3+ components takes longer for planetary than for stellar systems

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